Tate Homology and Powered Flybys

Kevin Ruck

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The Main Result

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The Main Result

Theorem 1

In the setting of the planar circular restricted three body problem

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In the setting of the planar circular restricted three body problem there are infinitely many symmetric consecutive collision orbits for all energies below the first critical energy value, which intersect their symmetry axis on the straight line between the second and the main body.



Powered Flybys

The Manoeuvre



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Powered Flybys

The Rough Idea:

A rocket uses the conservation of momentum to gain velocity



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Powered Flybys

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- $\rightarrow\,$ The gain of kinetic energy is big, if the starting velocity is high.
- \rightarrow This phenomenon is called the Oberth effect after Hermann Oberth (1923)

Powered Flybys

What about conservation of energy?

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Powered Flybys

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As the rocket gains kinetic energy, so does the fuel. If the fuel is burned the rocket profits from both the internal and the kinetic energy .

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Why consecutive collision orbits?

• For every c.c. orbit there is a flyby orbit close by.

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As the rocket gains kinetic energy, so does the fuel. If the fuel is burned the rocket profits from both the internal and the kinetic energy .

Why consecutive collision orbits?

- For every c.c. orbit there is a flyby orbit close by.
- We have a sharp mathematical definition for c.c. orbits.



Orientation of C.C. Orbits

Levi-Civita Regularization

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For this regularization procedure we interpret our phase space as \mathbb{C}^2 and pull back with the map

$$\mathcal{L}: \mathbb{C}^{2} \to \mathbb{C}^{2}; \quad (z, w) \mapsto \left(z^{2}, \frac{w}{2\overline{z}}\right)$$

$$H \xrightarrow{\text{regularise}} |z|^{2} \mathcal{L}^{*} H$$

$$(w, h_{\mathcal{C}}, \psi, -\mathcal{L}, \psi, -\mathcal{L}, \psi, \psi, -\mathcal{L},$$

Orientation of C.C. Orbits

Why Orientation matters:



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A Short Introduction to Lagrangian RFH

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A Short Introduction to Lagrangian RFH

Given a Hamiltonian system (M, ω, H) .

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$$\mathcal{A}_{H}(x, au) := \int\limits_{0}^{1} \lambda^{*} x \mathrm{d}t - au \int\limits_{0}^{1} H(x(t)) \mathrm{d}t + \mathrm{constant}(x)$$

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on a manifold $P(M, L) \times \mathbb{R}$, where

 $P(M,L) := \{x \text{ path in } M \text{ with } x(0) \in L, x(1) \in L\}$

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A Short Introduction to Lagrangian RFH

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The Lagrangian Rabinowitz Floer Homology is then defined as

$$RFH(M, H, L) := \frac{\ker \partial}{\operatorname{im} \partial}.$$

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$$dim RFH(M, H, L) \le \#crit A_H$$

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The powerfull property of RFH is that it only depends on the Hamiltonian system up to homotopy.

A Short Introduction to Lagrangian RFH

Problem:

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A Short Introduction to Lagrangian RFH

Problem: RFH = 0 in our setting

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A Short Introduction to Lagrangian RFH

Problem: RFH = 0 in our setting Solution:

A Short Introduction to Lagrangian RFH

Problem: RFH = 0 in our setting Solution: equivariant version of Lagrangian RFH

G-Equivariant Lagrangian RFH

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$$CRF \xrightarrow{/G} CRF^G$$
 $\partial \xrightarrow{/G} \partial^G$

G-Equivariant Lagrangian RFH

The idea of equivariant RFH is that dividing out a symmetry will improve the strength of the resulting homology.

$$CRF \xrightarrow{/G} CRF \subseteq \neg finite group, acts free
\partial \xrightarrow{/G} \partial^G_{\leftarrow} count only up to Groction
\Rightarrow RFH^G = \frac{\ker \partial^G}{2}$$

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Then the G-equivariant Lagrangian RFH is equal to the Tate homology of G (with \mathbb{Z}_2 coefficients), i.e.

$$RFH^G_*(M, H, L) = TH_*(G, \mathbb{Z}_2).$$

The Main Result

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- **2** Use the fact that L_{col} can be mapped to L_S via a Hamiltonian diffeomorphism to show that dim (RFH(M, H, L)) is a lower bound for the number of trajectories.

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So Remember that $TH_i(\mathbb{Z}_2, \mathbb{Z}_2) = \mathbb{Z}_2 \ \forall i \in \mathbb{Z}$