

# Floer-theoretic corrections to the geometry of moduli spaces of Lagrangian tori

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# SYZ mirror symmetry

Strominger-Yau-Zaslow (SYZ) construction of mirror to  $(X, \omega)$  (Kähler)  
 rel. anticanonical divisor  $D \subset X$  ( $[D] = c_1(TX)$ ) (log CY pair)

① find a **Lagrangian torus fibration**  $\pi: X - D \rightarrow B$  (w/ fibers of vanishing Maslov class,  
 generically unobstructed, ...)

→ uncorrected mirror of  $X - D$ :  $\left\{ (F_b, \nabla) \middle| \begin{array}{l} F_b = \pi^{-1}(b) \text{ fiber of } \pi \\ \nabla \text{ unitary rk. 1 loc. system/F} \end{array} \right\}$   
 analytic space, loc.  $\simeq (IK^*)^n$  (moduli space of Lagr. tori +  $\nabla$  in  $X - D$ ).

② Floer-theoretic corrections to geometry of  $X^\vee$  from **holomorphic discs** in  $(X, F_b)$

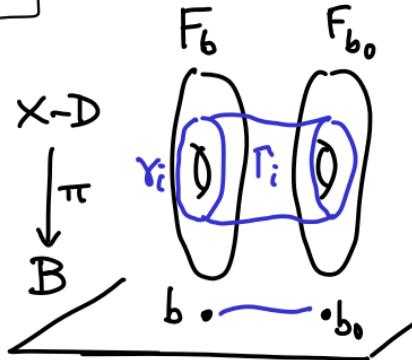
- Maslov index  $\mu(u) = 2[D] \cdot [u] = 0$  discs deform analytic structure of  $X^\vee$   
 (fix issues w/ singular fibers of  $\pi$ )

- $\mu=2$  discs ( $[D] \cdot [u] = 1$ ) → superpotential  $W \in \mathcal{O}(X^\vee)$   
 $(X^\vee, W)$  Landau-Ginzburg model.

NEW PHENOMENON

- $\mu < 0$  discs ( $[D] \cdot [u] < 0$ ) → extended deformations -  $X^\vee$  no longer analytic space!!

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## The uncorrected moduli space $X_0^\vee$

$\pi: X - D \rightarrow B$  Lagr. torus fibration

$$\rightarrow X_0^\vee := \left\{ (F_b, \nabla) \mid \begin{array}{l} F_b = \pi^{-1}(b) \text{ smooth fiber of } \pi \\ \nabla \in \text{hom}\left(\pi_1(F_b), U(1)_K\right) = H^1(F_b, U(1)_K) \end{array} \right\}$$

analytic space /  $K$  (away from sing. fibers):

$\{r_i\}$  basis of  $H_1(F_b)$   $\rightsquigarrow$  local coords.  $z_i(F_b, \nabla) = T^{\omega(r_i)} \cdot \nabla(r_i) \in K^*$ .  
+ base point  $b_0$

Here  $K = \left\{ \sum a_i T^{\lambda_i} \mid a_i \in k, \begin{array}{l} \lambda_i \in \mathbb{R} \\ \lambda_i \rightarrow +\infty \end{array} \right\}$  Novikov field in formal variable  $T$

- Floer-theoretic weights of holom. discs in  $(X, F_b)$ :  $F_b$    $[u] = \beta \in \pi_2(X, F_b)$ .

$z_\beta = T^{\omega(\beta)} \nabla(\partial\beta) \in K^*$  are monomials in  $z_1, \dots, z_n$

$\Rightarrow$  Lagr. Floer theory has analytic dependence on  $(F, \nabla) \in X_0^\vee$ .

$\rightsquigarrow$  family Floer approach to mirror symmetry (Fukaya, Abouzaid, Tu, Yuan, ...)

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## Holomorphic discs and the Floer-theoretic obstruction

If there are no holom. discs,  $X_0^V := \{(F_b, \nabla) \mid F_b = \pi^{-1}(b), \nabla \in \text{hom}(\pi_1(F_b), U(1)_K)\}$   
 is a moduli space of objects of  $\mathcal{F}(X \setminus D)$ , with its natural analytic structure.

Holomorphic discs deform the Fukaya category. **Floer obstruction** (FOOO)  $m_0^L \in CF(L, L)$

$$\mathcal{M}_1(X, L; J, \beta) = \left\{ u: (\mathbb{D}^2, \partial) \rightarrow (X, L) \mid \begin{array}{l} \bar{\partial}_J u = 0 \\ [u] = \beta \end{array} \right\} / \text{Aut}(\mathbb{D}^2, \partial) \xrightarrow{ev_\beta} L \quad \text{Diagram: } \mathbb{D}^2 \xrightarrow{u} X \text{ (with boundary point } \beta \text{ marked)}.$$

expected  $\dim_K = n - 2 + \mu(\beta)$ .

$$m_0^{(F_b, \nabla)} = \sum_{\beta \in \pi_2(X, F_b)} ev_{\beta_X} [\mathcal{M}_1(X, F_b; J, \beta)] \underbrace{T^{\omega(\beta)}}_{z_\beta} \nabla(\beta) \in C^*(F_b; K).$$

$\hookrightarrow$  stable map compactif.  $(\deg = 2 - \mu(\beta))$ .

Expect:  $\rightarrow \mu = 0$  discs (e.g. discs in  $X \setminus D$ ) occur along union of (thickened) codim\_K 1 walls in  $B$

$\rightarrow$  outside of walls,  $F_b$  is weakly unobstructed: min  $\mu = 2$ , so  $m_0 = W_{(F, \nabla)} \cdot 1$

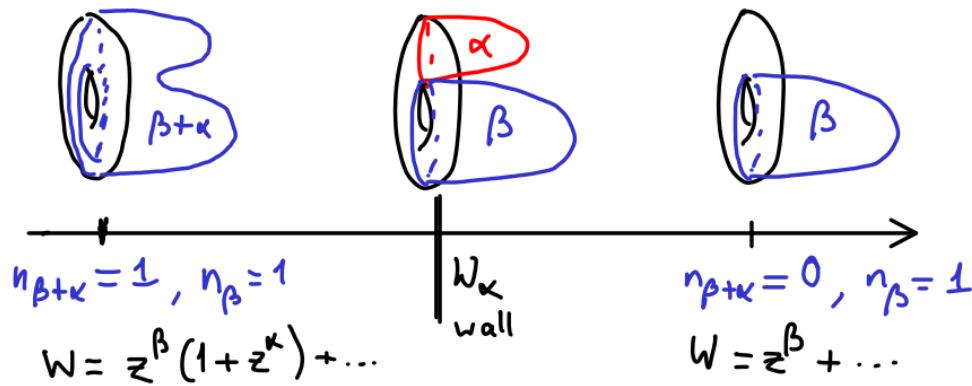
superpotential  $W(F_b, \nabla) = \sum_{\mu(\beta)=2} n_\beta z_\beta \in \mathcal{O}(X_0^V), n_\beta = \deg(ev_\beta) \in \mathbb{Z}.$

$\downarrow$  Floer cohomology still well-defined.

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## Wall-crossing

The superpotential  $W(F_b, \nabla) = \sum_{\mu(\beta)=2} n_\beta z_\beta$  is analytic over domains delimited by the walls in  $B$ , but counts of  $\mu=2$  discs have wall-crossing **discontinuities** due to bubbling of  $\mu=0$  discs.



But: formulas for  $W$  match under **wall-crossing coord. changes** ( $\varphi: z^\beta \mapsto z^\beta (1+z^\alpha + \dots)^{[\partial/\partial z] \cdot [W_\alpha]}$ )

→ **Corrected mirror**  $X^\nu =$  regime local pieces of  $X^{v_0}$  via these coord. changes;  $W \in \mathcal{O}(X^\nu)$  globally.

Consistency of wall-crossing transformations (cocycle condition) ⇒  $X^\nu$  well-defined analytic space

Kontsevich-Soibelman algorithm to construct consistent scattering diagram of walls (Gross-Siebert)

... but consistency fails when there are  $\mu < 0$  discs !!!

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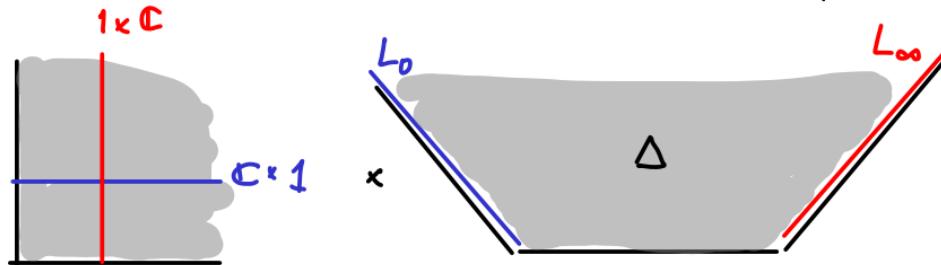
# Main example

$X = \text{blowup of toric 4-fold } \mathbb{C}^2 \times K_{\mathbb{P}^1} \text{ along } \mathbb{C} \times \{1\} \times L_0 \cup \{1\} \times \mathbb{C} \times L_\infty$

$D = \text{proper transform of } U \text{ toric strata}$

$\uparrow$   
 $\mathcal{O}_{\mathbb{P}^1}(-2)$

$\uparrow$   
 fibers of  $K_{\mathbb{P}^1} \rightarrow \mathbb{P}^1$  over 0 and  $\infty$

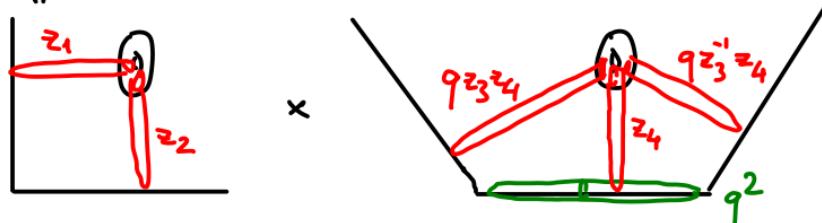


$$\pi: X \setminus D \rightarrow B \cong \mathbb{R}_+^2 \times \text{int}(\Delta)$$

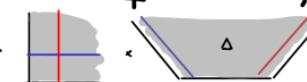


fibers away from exc. divisors  $E_0 \cup E_\infty = \text{lifts of product tori in } \mathbb{C}^2 \times K_{\mathbb{P}^1}$

For  $\mathbb{C}^2 \times K_{\mathbb{P}^1}$ , the mirror is  $(K^*)^4$ ,  $W = z_1 + z_2 + (1 + q^2 + q z_3 + q z_3^{-1}) z_4$



Toric  $\mu=2$  discs along coord axes, intersecting one toric divisor  
 $+ q^2 z_4$  term: disc  $\cup S \hookrightarrow \overset{\text{zero section in } K_{\mathbb{P}^1}}{N = \mathcal{O} \oplus \mathcal{O} \oplus \mathcal{O}(-2)}$

6]  $X = \text{Bl}(\mathbb{C}^2 \times K_{\mathbb{P}^1}, \mathbb{C} \times 1 \times L_0 \cup 1 \times \mathbb{C} \times L_\infty)$   $\pi: X \setminus D \rightarrow B \cong \mathbb{R}_+^2 \times \text{int}(\Delta)$   
 w/ singular fibers over 

fibers away from exc. divisors  $E_0 \cup E_\infty =$  lifts of product tori in  $\mathbb{C}^2 \times K_{\mathbb{P}^1}$

wall-crossing occurs as radii  $r_1, r_2$  in  $\mathbb{C}^2$   $\nearrow$  through 1: new  $\mu=2$  discs in  $X$ :

proper transforms of  $\mu \geq 4$  discs in  $\mathbb{C}^2 \times K_{\mathbb{P}^1}$  through  $\mathbb{C} \times 1 \times L_0$  ( $r_2 > 1$ ) or  $1 \times \mathbb{C} \times L_\infty$  ( $r_1 > 1$ ).

$(x_2, x_4)$  disc st.  $x_2 = 1 @ x_4 = 0$   
 union  $\widehat{\{(x_2, x_4) | x_2 = 1\}}$

$r_1 < 1$        $r_2 > 1$

$(x_1, x_2, x_4)$  disc st.  
 $(x_1, x_2) = (1, 1) @ x_4 = 0$

union  $\widehat{\{(x_1, x_2) = (1, 1)\}}$

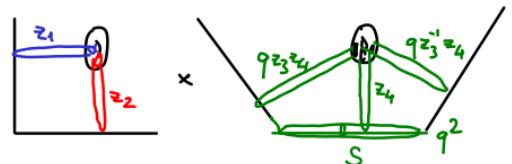
$$W_{-+} = z_1 + z_2 (1 + q q' z_4 + q' z_3 z_4) + (1 + q^2 + q z_3 + q z_3^{-1}) z_4$$

$r_2 > 1$

$$\begin{aligned} W_{++} = & z_1 (1 + q q'' z_4 + q'' z_3^{-1} z_4) \\ & + z_2 (1 + q q' z_4 + q' z_3 z_4) + q' q'' z_1 z_2 z_4 \\ & + (1 + q^2 + q z_3 + q z_3^{-1}) z_4 \end{aligned}$$

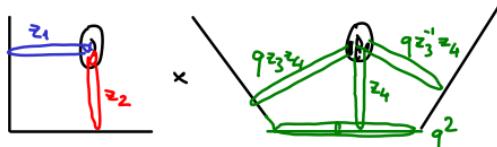
$$W_{--} = z_1 + z_2 + (1 + q^2 + q z_3 + q z_3^{-1}) z_4$$

same as for  $\mathbb{C}^2 \times K_{\mathbb{P}^1}$   
 (no discs through  $E_0 \cup E_\infty$ )



$$\begin{aligned} W_{+-} = & z_1 (1 + q q'' z_4 + q'' z_3^{-1} z_4) + z_2 \\ & + (1 + q^2 + q z_3 + q z_3^{-1}) z_4 \end{aligned}$$

7) The mirror of  $X = \text{Bl}(\mathbb{C}^2 \times K_{\mathbb{P}^1}, \mathbb{C} \times 1 \times L_0 \cup 1 \times \mathbb{C} \times L_\infty)$



Wall crossing at  $r_1=1$  preserves  $\bar{z}_2, \bar{z}_3, \bar{z}_4$   
 $r_2=1$  preserves  $\bar{z}_1, \bar{z}_3, \bar{z}_4$

$$\varphi_{0+}: z_1 \mapsto z_1(1 + qq''z_4 + q''\bar{z}_3\bar{z}_4 + q'q''z_2z_4)$$

$$r_2 > 1$$

$$W_{-+} = z_1 + z_2(1 + qq'z_4 + q'\bar{z}_3\bar{z}_4) + (1 + q^2 + qz_3 + qz_3^{-1})z_4$$

$$W_{++} = z_1(1 + qq''z_4 + q''\bar{z}_3\bar{z}_4) + z_2(1 + qq'z_4 + q'\bar{z}_3\bar{z}_4) + q'q''z_1z_2z_4 + (1 + q^2 + qz_3 + qz_3^{-1})z_4$$

$$\varphi_{-0}: z_2 \mapsto z_2(1 + qq'z_4 + q'\bar{z}_3\bar{z}_4)$$

$$\varphi_{+0}: z_2 \mapsto z_2(1 + qq'z_4 + q'\bar{z}_3\bar{z}_4 + q'q''z_1z_4)$$

$$r_2 < 1$$

$$W_{--} = z_1 + z_2 + (1 + q^2 + qz_3 + qz_3^{-1})z_4$$

$$W_{+-} = z_1(1 + qq''z_4 + q''\bar{z}_3\bar{z}_4) + z_2 + (1 + q^2 + qz_3 + qz_3^{-1})z_4$$

$$r_1 < 1$$

$$\varphi_{0-}: z_1 \mapsto z_1(1 + qq''z_4 + q''\bar{z}_3\bar{z}_4) \quad r_1 > 1$$

$\varphi_{+0} \circ \varphi_{0-} \neq \varphi_{0+} \circ \varphi_{-0}$ : the wall-crossing diagram is inconsistent!!

This is caused by  $\mu=-2$  stable disc at  $(r_1, r_2) = (1, 1)$ :  $z_4$ -disc  $\cup \widehat{S}_{(x_1, x_2) = (1, 1)}$  (weight  $q'q''z_4$ ).

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## $X^\vee$ as a deformed Landau-Ginzburg model.

The mirror constructed above is consistent mod  $q'q''$ . The extra terms can be viewed as a deformation of (alg. geom. of)  $(X^\vee, W)$  by a class in

$$\mathrm{HH}^*(\mathrm{MF}(X^\vee, W)) = \mathrm{H}^*(X^\vee, (\lambda^* T_{X^\vee}, \sharp_{dW}))$$

determined by  $\mu = -2$  discs : leading term  $w^{(2)} \in \check{\mathcal{C}}^2(X^\vee, \lambda^2 T_{X^\vee})$  takes value  $w_{00}^{(2)} = q'q'' z_4 \partial_{\log z_1} \wedge \partial_{\log z_2}$  on overlap of coord. charts around  $(r_1, r_2) = (1, 1)$ .

$$\left( \begin{array}{l} \text{complete to a cocycle by adding } w_{+0}^{(1)} = q'q'' z_1 z_4 \partial_{\log z_2} \\ \text{ } q'q'' \text{ terms in } \varphi_{10}, \varphi_{0+}, \varphi_{++} : \quad w_{0+}^{(1)} = q'q'' z_2 z_4 \partial_{\log z_1} \\ \quad w_{++}^{(0)} = q'q'' z_1 z_2 z_4 \end{array} \right) \Rightarrow \begin{array}{l} \sharp_{dW}(w^{(2)}) = \delta w^{(1)} \\ \sharp_{dW}(w^{(1)}) = \delta w^{(0)} \end{array}$$

$\rightarrow X^\vee$  deformed Landau-Ginzburg model:  $W$  loc. analytic, but coord. charts fail to satisfy cocycle condition by  $\sharp_{dW}(w^{(2)})$ .

9) A family Fiber approach to correcting the mirror

$\pi: X \sim D \rightarrow B$  Lagrangian fibration,  $X^{v0} = \{(F_b, \nabla) \mid F_b = \pi^{-1}(b) \text{ smooth fiber}\}$

$$\downarrow \pi^v \qquad \qquad \qquad \nabla \in \text{hom}(\pi_1(F_b), U(1)_K)$$

$\pi_*^v \mathcal{O} =: \mathcal{O}_{an}$  sheaf on  $B^\circ$   $B^\circ$  (smooth part) uncorrected mirror

= completion of  $K[H_1(F_b)] = K[z_1^{\pm 1}, \dots, z_n^{\pm 1}]$  ( $\Rightarrow$  monomials  $z^\beta$  = disc weights)

Lage. Floer theory for discs in  $(X, F_b)$   $\forall b \in B^\circ$ , with universal weights  $\in \mathcal{O}_{an}$

$\Rightarrow$  curved A<sub>∞</sub>-operations on  $C^*(B^\circ; C^*(F_b) \otimes \mathcal{O}_{an})$ .

$$m_0 = \sum_{\beta \in \pi_2(X, F_b)} ev_{\beta_X} [\bigcup_b \bar{M}_1(X, F_b; J, \beta)] z^\beta \in C^*(B^\circ; C^*(F_b) \otimes \mathcal{O}_{an}).$$

{local system of abelian groups on  $B^\circ$ }  $(\deg = 2 - \mu(\beta))$ .

Expect: || 1)  $\exists \alpha^{(i)} \in C^*(B^\circ; H^i(F_b) \otimes \mathcal{O}_{an})$  st.  $m_0 = \sum_{\mu=2} \alpha^{(0)} + \sum_{\mu=0} \alpha^{(1)} + \sum_{\mu=-2} \alpha^{(2)} + \dots$

Fukaya,  
K.Irie  
+ ...

|| 2)  $m_0$  satisfies master equation  $\delta m_0 + \frac{1}{2} \{m_0, m_0\}_{\text{deg. } -1 \text{ bracket}} = 0$

# The master equation & mirror geometry

- $m_0 = \alpha^{(0)} + \alpha^{(1)} + \dots \in \bigoplus_{i \geq 0} C^i(B^\circ; H^i(F_b) \hat{\otimes} \mathcal{O}_{\text{an}})$        $\alpha^{(i)}$  records codim.  $i$  walks  
of  $\mu = 2 - 2i$  discs
- on  $H^*(F_b) \otimes \mathbb{K}[H_1(F_b)] : \left\{ z^\gamma \alpha, z^{\gamma'} \alpha' \right\} := z^{\gamma+\gamma'} (\alpha \wedge {}_{\gamma'} \alpha' + (-1)^{|\alpha|} {}_{\gamma'} \alpha \wedge \alpha')$  (deg. -1)  
 $(\simeq H_*(\mathcal{L}F_b), \text{ Chas-Sullivan string bracket}) \quad \forall \alpha, \alpha' \in H^*(F_b) \quad \gamma, \gamma' \in H_1(F_b).$
- extend (w/ cup-product) to  $\{\cdot, \cdot\}$  on  $C^*(B^\circ; H^*(F_b) \hat{\otimes} \mathcal{O}_{\text{an}})$ .
- Using  $H^*(F_b, \mathbb{R}) \simeq T_b B$ , get  $H^*(F_b) \hat{\otimes} \mathcal{O}_{\text{an}} \xrightarrow{\sim} \pi_*^v(\Lambda^* T_{X^v 0})$   
 $z^\gamma r_{j_1}^* \wedge \dots \wedge r_{j_k}^* \longmapsto z^\gamma \partial_{\log z_{j_1}} \wedge \dots \wedge \partial_{\log z_{j_k}}$
- under which  $m_0 = \alpha^{(0)} + \alpha^{(1)} + \dots \longmapsto W = W^{(0)} + W^{(1)} + \dots \in \bigoplus_{i \geq 0} C^i(X^v 0, \Lambda^i T_{X^v 0})$   
 $\{\cdot, \cdot\} \longmapsto \text{Schouten-Nijenhuis bracket}$
- $\delta m_0 + \frac{1}{2} \{m_0, m_0\} = 0 \longmapsto \delta W + \frac{1}{2} [W, W] = 0$

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# The geometry of the master equation

$$\mathbb{W} = W^{(0)} + W^{(1)} + \dots \in \bigoplus_{i \geq 0} C^i(X^{\nu_0}, \Lambda^i T_{X^{\nu_0}}) \quad \text{Schouten-Nijenhuis bracket}$$

$$\delta \mathbb{W} + \frac{1}{2} [\mathbb{W}, \mathbb{W}] = 0 \iff (\delta + [W, \cdot])^2 = 0 \iff \begin{cases} (\delta + [W^{(1)}, \cdot]) W^{(0)} = 0 \\ (\delta + [W^{(1)}, \cdot])^2 = [z_{dW^{(0)}}(W^{(2)})] \cdot \end{cases} \quad (1)$$

$$(\delta + [W^{(1)}, \cdot])^2 = [z_{dW^{(0)}}(W^{(2)})] \cdot \quad (2)$$

$$(\delta + [W^{(1)}, \cdot]) W^{(2)} = z_{dW^{(0)}}(W^{(3)}) \quad (3)$$

View  $\delta + [W^{(1)}, \cdot]$  as *deform-* of analytic structure

$\left\{ \begin{array}{l} \text{Čech: deform gluing of } U_i \leftrightarrow U_j \text{ by vector field } W_{ij}^{(1)} \\ \text{Dolbeault: deform } \bar{\partial} \text{ operator by } W^{(1)} \in \Omega^{0,1} \otimes T^{1,0} \end{array} \right.$

(1)  $\Rightarrow W^{(0)}$  is analytic wrt deformed structure

(2)  $\Rightarrow$  deformation actually fails to be an analytic space by  $z_{dW^{(0)}}(W^{(2)}) \in C^2(X^\nu, T_{X^\nu})$

$\left\{ \begin{array}{l} \text{Čech: failure of coord. changes to satisfy cocycle condition} \\ \text{Dolbeault: failure of integrability of the deformed cx. structure} = \text{Nijenhuis tensor} \end{array} \right.$

Note:  $\text{crit}(W^{(0)})$  is still honestly analytic, since  $z_{dW^{(0)}}(W^{(2)}) = 0$  along critical locus.

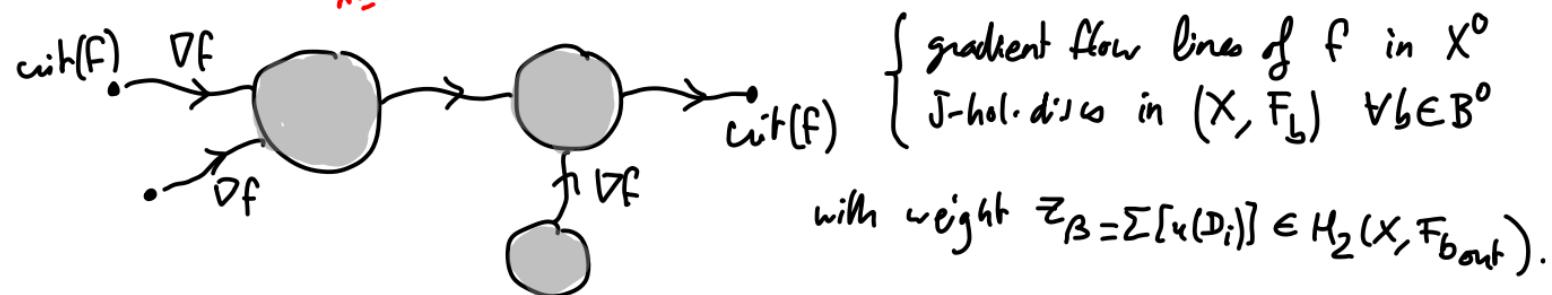
(12) A Morse model for Floer operations on  $C^*(B^0; H^*(F_b) \hat{\otimes} \mathcal{O}_{\text{an}})$

Fix a Morse function  $f: X^0 = \pi^{-1}(B^0) \rightarrow \mathbb{R}$  (+Morse-Smale metric)

triangulation  $\mathcal{P}$  of  $B^0$   $\rightsquigarrow$  Morse  $f|_{\mathcal{P}}$  on  $B^0$  with 1 cut pt  $b_\sigma$  index  $k$   $\forall \sigma \in \mathcal{P}^{[k]}$ ,  $\overline{w}(b_\sigma) = \sigma$ .  
+ add perfect Morse function on  $F_{b_\sigma}$ .

$\rightarrow CM^*(f) \simeq \check{C}^*(B^0; H^*(F_b))$  for open cover of  $B^0$  by stars of  $\text{Vert}(\mathcal{P})$

Define Floer operations  $(m_k)_{k \geq 0}$  on  $CM^*(f; \mathcal{O}_{\text{an}})$  by counting (perturbed) J-holom. treed discs



- achieve transversality using domain stabilization (cf. Charest-Woodward; divisors vary over  $B^0$ )  
+ domain-dependent perturbations
- convergence under assumption that  $\mathcal{P}$  fine enough ( $\Rightarrow \text{val}(z_{\beta_i}) > 0$  at  $b_{\text{out}}$ )