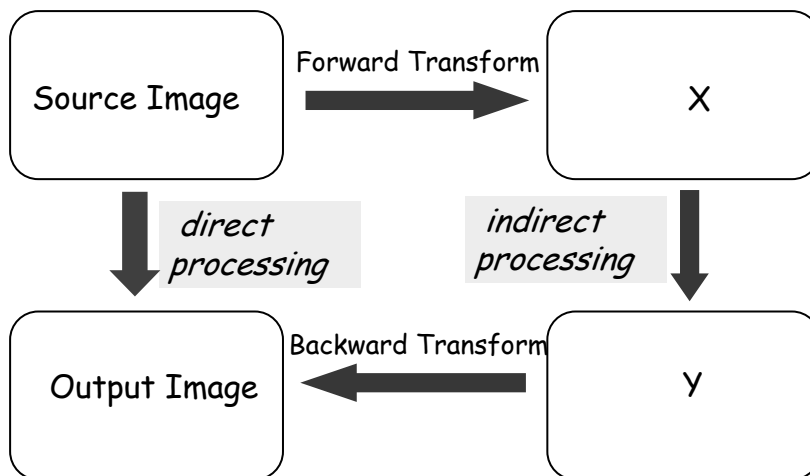


Image Transforms

- Fourier, FFT
- Wavelets
- Radon, Hough

A general note on image Transforms

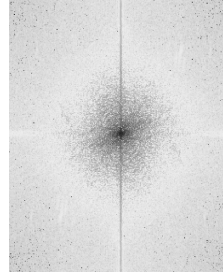
It is sometimes easier to a problem in a different space



Jean Baptiste Joseph Fourier

Fourier

FFT of Fourier

Image
domainFrequency
domain

$$f(x) = \sum_{n=-N}^{n=N} F_n e^{in\pi x / L}$$

Euler: $e^{i\theta} = \cos(\theta) + i \sin(\theta)$

Every digital signal
may be expanded in a
sum of sine & cosine
functions

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Simple Fourier expansions (Demo: uiFourier)

$$\text{Square} = \frac{\sin(x)}{1} + \frac{\sin(3x)}{3} + \frac{\sin(5x)}{5} + \frac{\sin(7x)}{7} + \frac{\sin(9x)}{9} + \dots$$

$$\text{Sawtooth} = \frac{\sin(x)}{1} + \frac{\sin(2x)}{2} + \frac{\sin(3x)}{3} + \frac{\sin(4x)}{4} + \frac{\sin(5x)}{5} + \dots$$

$$\text{Triangular} = \frac{\sin(x)}{1 \cdot 1} + \frac{\sin(3x)}{3 \cdot 3} + \frac{\sin(5x)}{5 \cdot 5} + \frac{\sin(7x)}{7 \cdot 7} + \frac{\sin(9x)}{9 \cdot 9} + \dots$$

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Fast Fourier Transform (FFT)

The function $f(x)$ is therefore completely specified by the expansion coefficients F_n

F is usually computed by an efficient algorithm discovered by Tukey and Cooley in 1965 - the FFT algorithm

FFT Demo (DMfft)

A simple rgb image compression using
fftn (multi dimensional fft)

```
% I = source image (rgb)
% mask = binary mask around the origin

FI = fftn(I);           % fft image
MFI = fftshift(mask).*FI; % masked fft image
RI = abs(ifftn(MFI));  % reconstruction
```

Fourier Transform

Infinite interval Case

$$X(\omega) = \int_{-\infty}^{\infty} x(t)e^{-i\omega t} dt$$

Inverse

$$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(\omega)e^{i\omega t} d\omega$$

Continuous Case - finite interval

$$X(\omega) = \sum_{n=-\infty}^{\infty} x_n e^{-in\omega}$$

Inverse

$$x_n = \frac{1}{2\pi} \int_{-\pi}^{\pi} X(\omega)e^{in\omega} d\omega$$

Discrete Version (DFT)

$$\hat{X}_v = \sum_{n=0}^{N-1} x_n e^{-\frac{2\pi i n v}{N}} \quad v = 0, 1, \dots, N-1$$

Inverse

$$x_n = \frac{1}{N} \sum_{v=0}^{N-1} \hat{X}_v e^{\frac{2\pi i n v}{N}} \quad n = 0, 1, \dots, N-1$$

Note: placement of constants is arbitrary. Only product of the constants in each direction counts.

Properties

Separability

$$\begin{aligned} F_{\mu\nu} &= \frac{1}{N^2} \sum_{n=0}^{N-1} \sum_{m=0}^{N-1} f_{nm} e^{-\frac{2\pi i(n\mu+m\nu)}{N}} & v, \nu = 0, 1, \dots, N-1 \\ &= \frac{1}{N} \sum_{n=0}^{N-1} F_{n\nu} e^{-\frac{2\pi i n \mu}{N}} \end{aligned}$$

where

$$F_{n\nu} = \frac{1}{N} \sum_{m=0}^{N-1} F_{nm} e^{-\frac{2\pi i m \nu}{N}}$$

Translation

$$F_{\mu-\mu_0} \leftrightarrow f_n e^{\frac{2\pi i n \mu_0}{N}}$$

$$F_{\mu} e^{-\frac{2\pi i n \mu_0}{N}} \leftrightarrow f_{n-n_0}$$

Convolution

Continuous

$$f * g(x) = \int f(y)g(x - y)dy$$

Discrete

$$f * g(x) = \frac{1}{N} \sum_{n=0}^{N-1} f(n)g(x - n)$$

Fourier

$$f(x, y) * g(x, y) \leftrightarrow F(u, v) \cdot G(u, v)$$

Correlation

Continuous

$$f \circ g(x) = \int f^*(y)g(x + y)dy = f^*(y) * g(-y)$$

Discrete

$$f \circ g(x) = \frac{1}{N} \sum_{n=0}^{N-1} f^*(n)g(x + n)$$

Fourier

$$f(x, y) \circ g(x, y) \leftrightarrow F^*(u, v) \cdot G(u, v)$$
$$f^*(x, y)g(x, y) \leftrightarrow F(u, v) \circ G(u, v)$$

Point Spread Function (PSF)

Response of a filter to a point image.

Let $P_{mn} = 1$ if $m = n = 0$ and 0 otherwise. Then

$$P'_{mn} = \sum_{i=-r}^r \sum_{j=-r}^r H_{ij}P_{m-i, n-j} = H_{mn}$$

Since convolution is linear when we know the effect of a filter on a point image we know it on any image.

Theorem: A linear shift-invariant operator must be a convolution operator in space.



Gaussian convolution

The solution of this equation ($D = 1$) is given by the **Convolution Integral**:

$$u(\vec{x}, t) = \begin{cases} f(\vec{x}) & (t = 0) \\ (G_{\sqrt{2t}} \otimes f)(\vec{x}) & (t > 0) \end{cases}$$

where

$$G_{\sigma}(\vec{x}) = \frac{1}{2\pi\sigma^2} \cdot \exp\left(-\frac{|\vec{x}|^2}{2\sigma^2}\right) \quad f: \mathbf{R}^2 \rightarrow \mathbf{R}$$



Johann Carl Friedrich Gauss



Born: 30 April 1777 in Brunswick, Duchy of Brunswick (now Germany)
Died: 23 Feb 1855 in Göttingen, Hanover (now Germany)

Heisenberg Uncertainty Principle (band limited in physical space and Fourier space)

Normalize $f(x)$ so that $\int_{-\infty}^{+\infty} |f(x)|^2 dx = 1$ and $\lim_{x \rightarrow \infty} \sqrt{x} f(x) = 0$

Let $F(k)$ be the Fourier transform of $f(x)$.

Then by Parseval's theorem $\int_{-\infty}^{+\infty} |F(k)|^2 dk = 1$

$$\begin{aligned}
 \left| \int_{-\infty}^{\infty} x(f(x)f'(x))dx \right| &\leq \sqrt{\int_{-\infty}^{\infty} |xf(x)|^2 dx} \sqrt{\int_{-\infty}^{\infty} |f'(x)|^2 dx} && \text{by the Schwarz inequality} \\
 &= \sqrt{\int_{-\infty}^{\infty} |xf(x)|^2 dx} \sqrt{\int_{-\infty}^{\infty} |ikF(k)|^2 dk} && \text{by Parseval's equality} \\
 \int_{-\infty}^{\infty} x(f(x)f'(x))dx &= \frac{1}{2} x [f(x)]^2 \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} \frac{1}{2} [f(x)]^2 \\
 &= 0 - \frac{1}{2} = -\frac{1}{2} && \text{by integration by parts}
 \end{aligned}$$

Hence

$$\sqrt{\int_{-\infty}^{\infty} |xf(x)|^2 dx} \sqrt{\int_{-\infty}^{\infty} |ikF(k)|^2 dk} \geq \frac{1}{2}$$

So we can't make both integrals small. If we concentrate the physical function we spread the Fourier transform and vice versa

We get equality if $f'(x) = \text{const } xf(x)$ or $f(x) = e^{-\alpha x^2}$. Hence, the Gaussian is the most concentrated function.

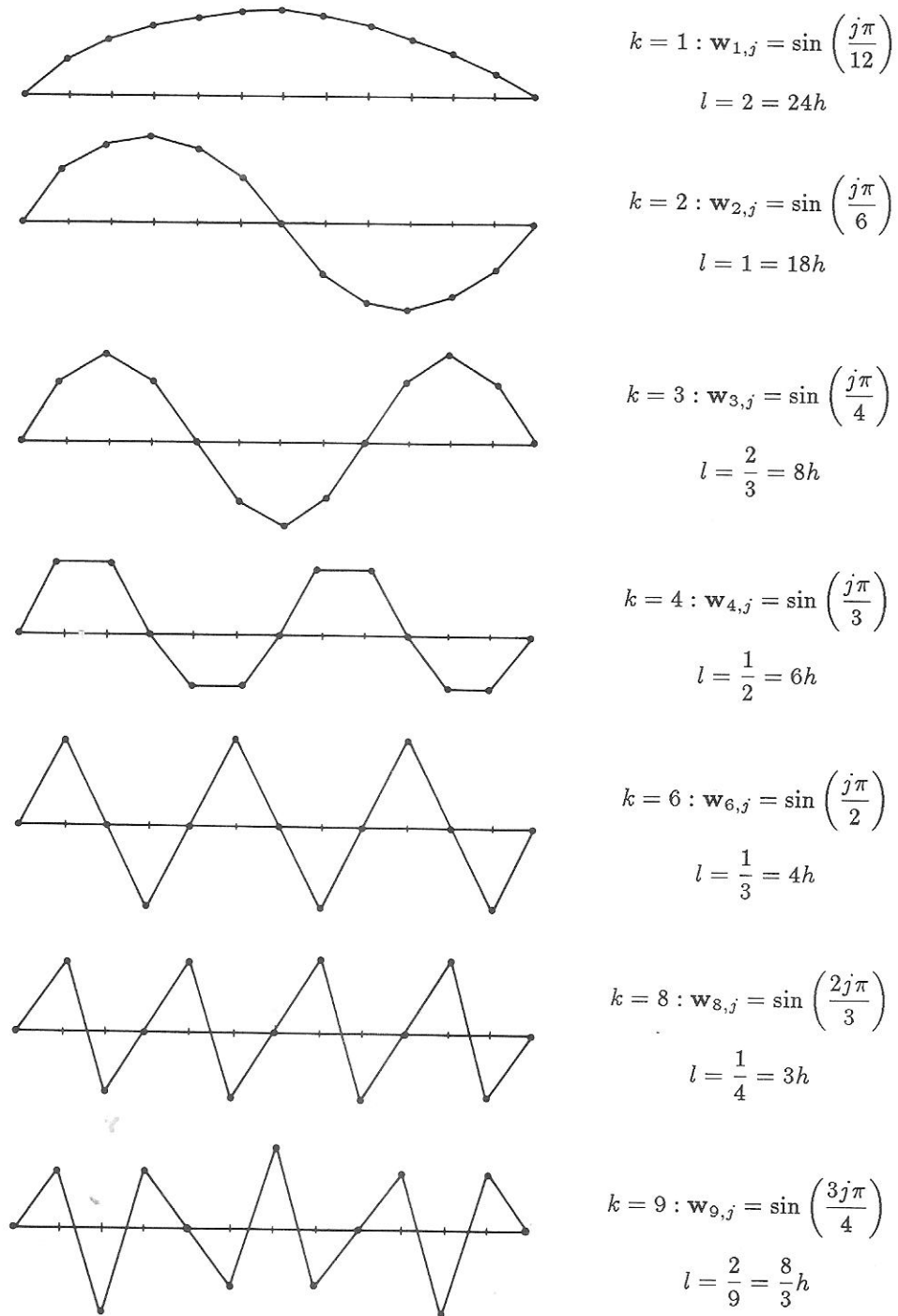


FIG. 8. Graphs of the Fourier modes of A on a grid with $N = 12$. Modes with wavenumbers $k = 1, 2, 3, 4, 6, 8, 9$ are shown. The wavelength of the k th mode is $l = 24h/k$.

$$w_{25-k,j} = (-1)^{j+1} w_{k,j}$$

FIG. 9.
eigenvalues
on the
 $1 \leq k \leq$

will on
error
H
smooth
the be
 $k \leq N$

Solving
for all
are re
damp
of an
schen
multi
ation
grid