

Fourier Series in Complex notation

$$\begin{aligned}\sin(x) &= \frac{e^{ix} - e^{-ix}}{2i} & \cos(x) &= \frac{e^{ix} + e^{-ix}}{2} \\ &= -\frac{i}{2} (e^{ix} - e^{-ix})\end{aligned}$$

So

$$\begin{aligned}\varphi(x) &= \frac{A_0}{2} + \sum_{n=1}^{\infty} A_n \cos\left(\frac{n\pi x}{l}\right) + B_n \sin\left(\frac{n\pi x}{l}\right) & -l < x < l \\ &= \frac{A_0}{2} + \sum_{n=1}^{\infty} A_n \frac{e^{in\pi x/l} + e^{-in\pi x/l}}{2} - \frac{i}{2} B_n (e^{in\pi x/l} - e^{-in\pi x/l}) \\ &= \frac{A_0}{2} + \sum_{n=1}^{\infty} \frac{A_n - iB_n}{2} e^{in\pi x/l} + \frac{A_n + iB_n}{2} e^{-in\pi x/l} \\ &= \sum_{n=-\infty}^{\infty} C_n e^{\frac{in\pi x}{l}} & C_n &= \begin{cases} \frac{A_n - iB_n}{2} & n > 0 \\ \frac{A_{-n} + iB_{-n}}{2} & n < 0 \\ \frac{A_0}{2} & n = 0 \end{cases}\end{aligned}$$

So $C_{-n} = C_n^*$

$$\begin{aligned}\int_{-l}^l e^{\frac{in\pi x}{l}} e^{-\frac{im\pi x}{l}} dx &= \begin{cases} 0 & m \neq n \\ 2l & m = n \end{cases} \\ C_n &= \frac{1}{2l} \int_{-l}^l \varphi(x) e^{-\frac{in\pi x}{l}} dx\end{aligned}$$

Sturm-Liouville

$$\begin{aligned}
 (*) \quad & -\frac{d}{dx} \left(p(x) \frac{du}{dx} \right) + q(x)u(x) = \lambda m(x)u(x) \quad a < x < b \\
 & p(x) > 0 \quad q(x) \geq 0 \quad m(x) > 0
 \end{aligned}$$

and all quantities are real.

Definition 1 λ is the eigenvalue and u is the eigenfunction

Definition 2 A homogeneous boundary condition is **symmetric** if

$$p[f'g' - f'g]_a^b = 0$$

Examples:

- Dirichlet
- Neumann
- periodic
- Robin (?)

$$\begin{aligned}
 a_l u(a) + b_l \frac{du}{dx}(a) &= 0 \\
 a_r u(b) + b_r \frac{du}{dx}(b) &= 0
 \end{aligned}$$

then

$$p[f'g' - f'g]_a^b = f(a) \left[-\frac{a_l}{b_l} g(a) \right] - \left(-\frac{a_l}{b_l} \right) f(a)g(a) = 0$$

Definition 3 Inner product

$$(u, v) = \int_a^b m(x)u(x)v(x)dx \quad m > 0$$

Example

$$p = 1 \quad q = 0 \quad m = 1$$

Then with Dirichlet conditions we have

$$\begin{aligned} -\frac{d^2u}{dx^2} &= \lambda u(x) & 0 < x < l \\ u(0) &= u(l) = 0 \end{aligned}$$

If $\lambda > 0$ then

$$u(x) = A \cos(\sqrt{\lambda}x) + B \sin(\sqrt{\lambda}x)$$

Using $u(0) = 0$ we get

$$u(x) = B \sin(\sqrt{\lambda}x)$$

and using $u(l) = 0$

$$\begin{aligned} \lambda &= \left(\frac{n\pi}{l}\right)^2 \\ u(x) &= B \sin\left(\frac{n\pi}{l}x\right) \end{aligned}$$

So we have an infinite number of eigenvalues/eigenfunctions.

Green's Identities

First Identity:

$$\int_a^b \frac{d}{dx} \left(p(x) \frac{du}{dx} \right) v(x) dx + \int_a^b p(x) \frac{du}{dx} \frac{dv}{dx} dx = \left[p(x) \frac{du}{dx} v \right]_a^b$$

Second Identity:

$$\int_a^b -\frac{d}{dx} \left(p(x) \frac{du}{dx} \right) v(x) dx + \int_a^b \frac{d}{dx} \left(p(x) \frac{dv}{dx} \right) u(x) dx = \left[p(x) \left(-\frac{du}{dx} v + \frac{dv}{dx} u \right) \right]_a^b$$

In multi-dimensions this generalizes to ($p = 1$)

$$\begin{aligned} \iiint_D (\nabla u \cdot \nabla v + v \Delta u) dV &= \iint_{\partial D} v \frac{\partial u}{\partial n} dS \\ \iiint_D (v \Delta u - u \Delta v) dV &= \iint_{\partial D} \left(v \frac{\partial u}{\partial n} - u \frac{\partial v}{\partial n} \right) dS \end{aligned}$$

Proof. From the divergence theorem

$$\begin{aligned}\iiint \operatorname{div}(F) dV &= \iint F \cdot n dS \\ \operatorname{div}(v \operatorname{grad} u) &= \nabla u \cdot \nabla v + v \Delta u \\ \text{So } \iiint (\nabla u \cdot \nabla v + v \Delta u) dV &= \iiint \operatorname{div}(v \operatorname{grad} u) dV \\ &= \iint v \operatorname{grad} u \cdot n dS = \iint v \frac{\partial u}{\partial n} dS\end{aligned}$$

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Interchange u and v and subtract to get the second identity.

From (*)

$$\begin{aligned} (1) \quad & -\frac{d}{dx} \left(p(x) \frac{du}{dx} \right) + q(x)u(x) = \lambda_1 m(x)u(x) \\ (2) \quad & -\frac{d}{dx} \left(p(x) \frac{dv}{dx} \right) + q(x)v(x) = \lambda_2 m(x)v(x) \end{aligned}$$

Multiply (1) by v and (2) by u and subtract and integrate

$$\int_a^b \left[-v(x) \frac{d}{dx} \left(p(x) \frac{du}{dx} \right) + u(x) \frac{d}{dx} \left(p(x) \frac{dv}{dx} \right) \right] dx = (\lambda_1 - \lambda_2) \int_a^b m(x)u(x)v(x)dx$$

Integrate by parts (Green's theorem in multidimensions)

$$\begin{aligned} \int_a^b \left[\frac{dv}{dx} \left(p(x) \frac{du}{dx} \right) - \frac{du}{dx} \left(p(x) \frac{dv}{dx} \right) \right] dx + \left[p \left(\frac{du}{dx} v - \frac{dv}{dx} u \right) \right]_a^b &= (\lambda_1 - \lambda_2) \int_a^b m(x)u(x)v(x)dx \\ \left[p \left(\frac{du}{dx} v - \frac{dv}{dx} u \right) \right]_a^b &= (\lambda_1 - \lambda_2) \int_a^b m(x)u(x)v(x)dx \end{aligned}$$

If the boundary conditions are symmetric then

$$(\lambda_1 - \lambda_2) \int_a^b m(x)u(x)v(x)dx = 0$$

Hence, if $\lambda_1 \neq \lambda_2$

Theorem 4 *For symmetric boundary conditions , if $\lambda_1 \neq \lambda_2$ then $(u, v) = 0$
If $\lambda_1 = \lambda_2$ then we have a subspace and we can choose an orthogonal basis. This
is done by Gram Schmidt*

Gram-Schmidt

If $\{\psi_k(x)\}$ is a linearly independent basis then we can construct an orthonormal basis that spans the same space.

- $\varphi_1(x) = \psi_1(x)$
- $\varphi_2(x) = \psi_2(x) - \frac{(\psi_2, \varphi_1)}{(\varphi_1, \varphi_1)} \varphi_1$
Then $(\varphi_2, \varphi_1) = (\psi_2, \varphi_1) - \frac{(\psi_2, \varphi_1)}{(\varphi_1, \varphi_1)} (\varphi_1, \varphi_1) = 0$
- $\varphi_k(x) = \psi_k(x) - \sum_{j=1}^{k-1} \frac{(\psi_k, \varphi_j)}{(\varphi_j, \varphi_j)} \varphi_j$

Hence, we can consider all the solutions of (*) to be orthogonal to each other. So consider a sequence of orthogonal solutions of (*) $\{\varphi_k(x)\}$. Then if

$$\begin{aligned} f(x) &= \sum_n a_n \varphi_n(x) \\ (f, \varphi_m) &= \sum_n a_n (\varphi_n, \varphi_m) \\ &= a_m (\varphi_m, \varphi_m) \end{aligned}$$

So

$$a_m = \frac{(f, \varphi_m)}{(\varphi_m, \varphi_m)} = \frac{\int_a^b m(x) f(x) \varphi_m(x) dx}{\int_a^b m(x) \varphi_m^2(x) dx}$$

Theorem 5 *If p, q, m are real and the boundary conditions are symmetric then there are no complex eigenvalues*

Proof.

$$\begin{aligned} (1) \quad & -\frac{d}{dx} \left(p(x) \frac{du}{dx} \right) + q(x)u(x) = \lambda m(x)u(x) \\ (2) \quad & -\frac{d}{dx} \left(p(x) \frac{d\bar{u}}{dx} \right) + q(x)\bar{u}(x) = \bar{\lambda} m(x)\bar{u}(x) \end{aligned}$$

As before multiply first equation by \bar{u} , the second by u , subtract and integrate. Then

$$\int_a^b \left[-\bar{u}(x) \frac{d}{dx} \left(p(x) \frac{du}{dx} \right) + u(x) \frac{d}{dx} \left(p(x) \frac{d\bar{u}}{dx} \right) \right] dx = (\lambda - \bar{\lambda}) \int_a^b m(x) u(x) \bar{u}(x) dx$$

Again integrate by parts and use the symmetry of the boundary conditions.

$$(\lambda - \bar{\lambda}) \int_a^b m(x) |u(x)|^2 dx = 0$$

So $\lambda = \bar{\lambda}$ i.e. λ is real

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If u is complex then its real and imaginary components are solutions.

Negative Eigenvalues

$$(*) \quad -\frac{d}{dx} \left(p(x) \frac{du}{dx} \right) + q(x)u(x) = \lambda m(x)u(x) \quad a < x < b$$

By Green's first identity we have for all v

$$\int_a^b \frac{d}{dx} \left(p(x) \frac{du}{dx} \right) v(x) dx = - \int_a^b p(x) \frac{du}{dx} \frac{dv}{dx} dx + \left[p(x) \frac{du}{dx} v \right]_a^b$$

Choose $v = u$ and assume symmetric boundary conditions. Then

$$\int_a^b u(x) \frac{d}{dx} \left(p(x) \frac{du}{dx} \right) dx = - \int_a^b p(x) \left(\frac{du}{dx} \right)^2 dx \leq 0$$

Using the ODE we get

$$\int_a^b u(x) [q(x)u(x) - \lambda m(x)u(x)] dx = - \int_a^b p(x) \left(\frac{du}{dx} \right)^2 dx$$

So

$$\begin{aligned} \lambda \int_a^b m(x)u^2(x)dx &= \int_a^b q(x)u^2(x)dx + \int_a^b p(x) \left(\frac{du}{dx} \right)^2 dx \\ \lambda &= \frac{\int_a^b q(x)u^2(x)dx + \int_a^b p(x) \left(\frac{du}{dx} \right)^2 dx}{\int_a^b m(x)u^2(x)dx} \geq 0 \end{aligned}$$

We can have equality only if $q(x) = 0$ and $\frac{du}{dx} = 0$

Note: All these proofs work equally well in multidimensions using Green's theorem instead of integration by parts.

Completeness

Theorem 6 *There are an infinite number of eigenvalues for (*) and $\lambda_n \rightarrow \infty$.
Furthermore*

$$f(x) = \sum c_n \varphi_n(x)$$

$$c_n = \frac{(f, \varphi_n)}{(\varphi_n, \varphi_n)}$$

Convergence:

$$\sum_{n=1}^{\infty} \varphi_n(x) \stackrel{?}{\rightarrow} f(x)$$

Definition 7 *Pointwise Convergence:*

$$\lim_{N \rightarrow \infty} \left| f(x) - \sum_{n=1}^N \varphi_n(x) \right| = 0 \quad \text{for every } x$$

Definition 8 *Uniform Convergence:*

$$\lim_{N \rightarrow \infty} \max_{a \leq x \leq b} \left| f(x) - \sum_{n=1}^N \varphi_n(x) \right| = 0 \quad \text{for every } x$$

Definition 9 L_2 (root mean square)

$$\lim_{N \rightarrow \infty} \int_a^b \left| f(x) - \sum_{n=1}^N \varphi_n(x) \right|^2 dx = 0$$

Uniform convergence implies pointwise convergence.

Uniform convergence implies root mean square convergence.

Examples

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$$f(x) = x^n \quad 0 \leq x \leq 1$$

Then

$$x_n \rightarrow \begin{cases} 0 & 0 < x < 1 \\ 1 & x = 1 \end{cases}$$

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$$f_n(x) = (1-x)x^{n-1} = x^{n-1} - x^n$$
$$\sum_{n=1}^N f_n(x) = \sum_{n=1}^N (x^{n-1} - x^n) = 1 - x^N \rightarrow 1 \quad \text{as } N \rightarrow \infty$$

so we have pointwise convergence.

However

$$\max_{0 \leq x \leq 1} |1 - (1 - x^N)| = \max_{0 \leq x \leq 1} |x^N| = 1 \neq 0$$

So we don't have uniform convergence.

For L_2 we have

$$\int_0^1 |x^N|^2 dx = \frac{1}{2N+1} \rightarrow 0$$

So we have L_2 convergence.

Theorem 10 *If*

- f, f', f'' exist and are continuous in $a \leq x \leq b$ i.e. $f \in C^2[a, b]$
- f satisfies the boundary conditions

then

$$f(x) = \sum_{n=1}^{\infty} a_n \varphi_n(x) \quad \text{converges uniformly}$$

Theorem 11 *If $\int_a^b f^2(x) dx < \infty$*

$$\text{then } f(x) = \sum_{n=1}^{\infty} a_n \varphi_n(x) \quad \text{converges in } L_2$$

Theorem 12 *For sine and cosine series only.*

If

- f is continuous on $a \leq x \leq b$
- f' is piecewise continuous

Then the series converges pointwise. If f and f' are piecewise continuous then

$$\sum a_n \varphi_n(x) \rightarrow \frac{f(x+) + f(x-)}{2}$$

Theorem 13 *Integration: If formally*

$$f(x) \leftrightarrow \frac{A_0}{2} + \sum_{n=1}^{\infty} A_n \cos\left(\frac{n\pi x}{l}\right) + B_n \sin\left(\frac{n\pi x}{l}\right) \quad \text{not necessarily convergent}$$

Then

$$\int_{-l}^x f(y) dy = \frac{A_0}{2}(x+l) + \frac{l}{\pi} \sum_{n=1}^{\infty} \left[\frac{A_n}{n} \sin\left(\frac{n\pi t}{l}\right) - \frac{B_n}{n} \cos\left(\frac{n\pi t}{l}\right) \right]_{-l}^x \quad \text{is convergent}$$

We note that for differentiation it is the opposite i.e. the derivative of a convergent series may not converge.

Example: expanding x in a sine series we have

$$x = 2 \sum_{n=1}^{\infty} \frac{l}{n\pi} (-1)^{n+1} \sin\left(\frac{n\pi x}{l}\right) \quad 0 \leq x \leq l$$

Differentiating we get

$$1 = 2 \sum_{n=1}^{\infty} (-1)^{n+1} \cos\left(\frac{n\pi x}{l}\right) \quad 0 \leq x \leq l$$

This is certainly **NOT** the cosine series of 1 which is just 1. In fact this series does not converge!

We begin with the proof of convergence in least squares. Restating the theorem we have

Theorem 14 If φ_n are the eigenfunctions of a Sturm-Liouville problem with symmetric boundary conditions and $\|f\| < \infty$. Then

$$\|f - \sum_{n=1}^N a_n \varphi_n\| \rightarrow 0 \quad a_n = \frac{(f, \varphi_n)}{(\varphi_n, \varphi_n)}$$

Theorem 15 Let φ_n be an orthogonal set and $\|f\| < \infty$, Then the choice of constants c_n that minimizes $\|f - \sum_{n \leq N} c_n \varphi_n\|$ is $c_n = a_n$

Proof. Assume for simplicity that all quantities are real

$$\begin{aligned} E_N &= \|f - \sum_{n \leq N} c_n \varphi_n\|^2 = \int |f(x) - \sum_{n \leq N} c_n \varphi_n(x)|^2 dx \\ &= \int |f(x)|^2 - 2 \sum_{n \leq N} c_n \int f(x) \varphi_n(x) dx + \sum_n \sum_m c_n c_m \int \varphi_n(x) \varphi_m(x) dx \\ &= \|f\|^2 - 2 \sum_{n \leq N} c_n (f, \varphi_n) + \sum_{n \leq N} c_n^2 \|\varphi_n(x)\|^2 \\ &= \|f\|^2 + \sum_{n \leq N} \|\varphi_n(x)\|^2 \left[c_n - \frac{(f, \varphi_n)}{\|\varphi_n(x)\|^2} \right]^2 - \sum_{n \leq N} \frac{(f, \varphi_n)^2}{\|\varphi_n(x)\|^2} \end{aligned}$$

To minimize we can only "play" with c_n . Since the middle term is positive we minimize E_N if

$$c_n = \frac{(f, \varphi_n)}{\|\varphi_n(x)\|^2} = a_n$$

Then

$$E_N = \|f\|^2 - \sum_{n \leq N} \frac{(f, \varphi_n)^2}{\|\varphi_n(x)\|^2} = \|f\|^2 - \sum_{n \leq N} a_n^2 \|\varphi_n(x)\|^2 \geq 0$$

So we have Bessel's inequality. If $\|f\|^2 < \infty$ then

$$\sum_{n \leq N} \frac{(f, \varphi_n)^2}{\|\varphi_n(x)\|^2} \leq \|f\|^2$$

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Parseval's Equality

Theorem 16 The Fourier Series converges to $f(x)$ in L_2 if and only if

$$\sum_{n \leq N} \frac{(f, \varphi_n)^2}{\|\varphi_n(x)\|^2} = \sum_{n \leq N} a_n^2 \|\varphi_n(x)\|^2 = \|f\|^2$$

Definition 17 A sequence $\{\varphi_n(x)\}$ is complete if Parseval's equality holds whenever $\|f\|^2 < \infty$

Riemann-Lebesgue Theorem

Theorem 18 If (a) $f \in C^1$
or (b) $\|f\|_{L^2} < \infty$

Then

$$\lim_{n \rightarrow \infty} \int_{-l}^l f(x) \begin{cases} \sin(\frac{n\pi x}{l}) \\ \cos(\frac{n\pi x}{l}) \end{cases} dx = 0$$

Proof.

1. integration by parts

2. In a Fourier series $B_n = \int_{-l}^l f(x) \sin(\frac{n\pi x}{l}) dx$ by Bessel's inequality $B_n \rightarrow 0$

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Example

Consider $f(x) = 1$ on $(0, \pi)$. We find that

$$1 = \sum_{n \text{ odd}} \frac{4}{n\pi} \sin(nx)$$

So by Parseval's equality

$$\begin{aligned} \int_0^\pi 1^2 dx &= \sum_{n \text{ odd}} \left(\frac{4}{n\pi} \right)^2 \frac{\pi}{2} \\ \sum_{n \text{ odd}} \frac{1}{n^2} &= \frac{\pi^2}{8} \end{aligned}$$

Pointwise Convergence

If

$$\begin{aligned} f(x) &= \frac{A_0}{2} + \sum_{n=1}^{\infty} A_n \cos(nx) + B_n \sin(nx) \\ A_n &= \frac{1}{\pi} \int_{-\pi}^{\pi} f(y) \cos(ny) dy \quad n = 0, 1, 2, 3... \\ B_n &= \frac{1}{\pi} \int_{-\pi}^{\pi} f(y) \sin(ny) dy \quad n = 1, 2, 3... \end{aligned}$$

Dirichlet kernel

Consider the partial sum

$$\begin{aligned}
S_N &= \frac{A_0}{2} + \sum_{n=1}^N A_n \cos(nx) + B_n \sin(nx) \\
&= \frac{1}{2\pi} \int_{-\pi}^{\pi} \left[1 + 2 \sum_{n=1}^N (\cos(nx) \cos(ny) + \sin(nx) \sin(ny)) \right] f(y) dy \\
&= \frac{1}{2\pi} \int_{-\pi}^{\pi} \left[1 + 2 \sum_{n=1}^N \cos(nx - ny) \right] f(y) dy \\
&= \frac{1}{2\pi} \int_{-\pi}^{\pi} K_N(x - y) f(y) dy
\end{aligned}$$

where

$$K_N(\theta) = 1 + 2 \sum_{n=1}^N \cos(n\theta) = \frac{\sin(N + \frac{1}{2})\theta}{\sin(\frac{\theta}{2})}$$

Proof. Use $\cos(\theta) = \frac{e^{in\theta} + e^{-in\theta}}{2}$ and get geometric series. ■

Now let $\theta = y - x$. Then

$$\begin{aligned}
S_N &= \frac{1}{2\pi} \int_{-\pi}^{\pi} K_N(\theta) f(x + \theta) d\theta \\
S_N(x) - f(x) &= \frac{1}{2\pi} \int_{-\pi}^{\pi} K_N(\theta) [f(x + \theta) - f(x)] d\theta \\
&= \frac{1}{2\pi} \int_{-\pi}^{\pi} g(\theta) \sin(N + \frac{1}{2})\theta d\theta \quad \text{where } g(\theta) = \frac{f(x + \theta) - f(x)}{\sin(\frac{\theta}{2})}
\end{aligned}$$

Let $\phi_n(\theta) = \sin(N + \frac{1}{2})\theta$. By Bessel's inequality we have

$$\sum_{n=1}^{\infty} \frac{|(g, \phi_n)|^2}{\|\phi_n\|^2} = \frac{1}{\pi} \sum_{n=1}^{\infty} |(g, \phi_n)| \leq \|g\|^2 = \int_{-\pi}^{\pi} \frac{[f(x + \theta) - f(x)]^2}{\sin^2(\frac{\theta}{2})} d\theta$$

By L'hospital's rule the integrand is finite at $\theta = 0$. Hence it is bounded everywhere and the integral exists. Since the sum converges each term must approach zero and so

$$|(g, \phi_n)| = \int_{-\pi}^{\pi} g(\theta) \sin(N + \frac{1}{2})\theta d\theta \rightarrow 0$$

Gibbs Phenomena

If instead we are interested in uniform convergence we need to analyze

$$\lim_{N \rightarrow \infty} \max_x |S_N(x) - f(x)|$$

One can show that if the function $f(x)$ has a discontinuity at $x = x_0$ then in fact this limit is nonzero and is about 9% of the size of the jump on either side.